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Joint determination of ⁴⁰K decay constants and ⁴⁰Ar*/⁴⁰K for the Fish Canyon sanidine standard, and improved accuracy for ⁴⁰Ar/³⁹Ar geochronology

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Abstract

⁴⁰Ar/³⁹Ar and K–Ar geochronology have long suffered from large systematic errors arising from imprecise K and Ar isotopic data for standards and imprecisely determined decay constants for the branched decay of ⁴⁰K by electron capture and β⁻ emission. This study presents a statistical optimization approach allowing constraints from ⁴⁰K activity data, K–Ar isotopic data, and pairs of ²³⁸U–²⁰⁶Pb and ⁴⁰Ar/³⁹Ar data for rigorously selected rocks to be used as inputs for estimating the partial decay constants (λ_e and λ_β) of ⁴⁰K and the ⁴⁰Ar^{*}/⁴⁰K ratio (κ_{FCs}) of the widely used Fish Canyon sanidine (FCs) standard. This yields values of $\kappa_{FCs} = (1.6418 \pm 0.0045) \times 10^{-3}$, $\lambda_e = (0.5755 \pm 0.0016) \times 10^{-10} a^{-1}$ and $\lambda_\beta = (4.9737 \pm 0.0093) \times 10^{-10} a^{-1}$. These results improve uncertainties in the decay constants by a factor of >4 relative to values derived from activity data alone. Uncertainties in these variables determined by our approach are moderately to highly correlated (cov(κ_{FCs}, λ_e) = 7.1889 × 10⁻¹⁹, cov($\kappa_{FCs}, \lambda_{\beta}$) = -7.1390 × 10⁻¹⁹, cov($\lambda_e, \lambda_{\beta}$) = -7.4497 × 10⁻²⁶) and one must take account of the covariances in error propagation by either linear or Monte Carlo methods. ⁴⁰Ar/³⁹Ar age errors estimated from these results are significantly reduced relative to previous calibrations. Also, age errors are smaller for a comparable level of isotopic measurement precision than those produced by the ²³⁸U/²⁰⁶Pb system, because the ⁴⁰Ar/³⁹Ar ages is now jointly calibrated by both the ⁴⁰K and ²³⁸U decay constants, and because $\lambda_e(^{40}K) < \lambda(^{238}U)$. Based on this new calibration, the age of the widely used Fish Canyon sanidine standard is 28.305 ± 0.036 Ma. The increased accuracy of ⁴⁰Ar/³⁹Ar ages is now adequate to provide meaningful validation of high-precision U/Pb or astronomical tuning ages in cases where closed system behavior of K and Ar can be established.

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1. INTRODUCTION

The ⁴⁰Ar/³⁹Ar dating method is one of the most important means of measuring geologic time over most of Earth's history. In addition to its wide age range of applicability

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and the existence of internal reliability criteria, this technique is of great importance because it is capable of exemplary precision, better than $0.1\%^1$ of the age in many cases, which is attainable by only the U/Pb method among widely used geochronometers. The 40 Ar/ 39 Ar method is especially important in the Cenozoic Era (the past 66 million years), and most currently used Cenozoic time scales (e.g., Cande

¹ All uncertainties herein are expressed at the level of one standard deviation or 68% confidence.

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and Kent, 1995; Aguilar et al., 1996; Gradstein et al., 2004) are calibrated almost exclusively by this method whereas U–Pb data are favored in the pre-Mesozoic. ${}^{40}Ar/{}^{39}Ar$ dating has been employed over the widest temporal range of any dating method, having illuminated important questions in geology, cosmology, archeology, and paleontology through the age of emplacement (i.e., volcanic) or cooling (i.e., plutonic and metamorphic) of rocks. ${}^{40}Ar/{}^{39}Ar$ dating is based on the ${}^{40}K-{}^{40}Ar$ method,

which requires knowledge of the decay constants both for the ${}^{40}\text{K} \rightarrow {}^{40}\text{Ar}$ electron capture decay (λ_{ϵ}) and the ${}^{40}\text{K} \rightarrow {}^{40}\text{Ca} \beta^- \text{ decay } (\lambda_\beta). A \beta^+ \text{ decay mode was inferred}$ by Beckinsale and Gale (1969) to account for 0.001% of ⁴⁰K activity, and has subsequently been ignored by other workers - this decay mode has never been detected to our knowledge. The ⁴⁰Ar/³⁹Ar method requires, in addition, the use of natural standards with either accurately-known ⁴⁰Ar/⁴⁰K ratios or independent knowledge of the age. The ⁴⁰K decay constants in widespread use since 1977 (Steiger and Jäger, 1977) are derived from a compilation of ⁴⁰K activity data (Beckinsale and Gale, 1969) combined with a revised value for the isotopic abundance of ⁴⁰K (Garner et al., 1975). Following the suggestion of Steiger and Jäger (1977) the IUGS Subcommission on Geochronology adopted a value for λ_{tot} (= $\lambda_{\varepsilon} + \lambda_{\beta}$) of 5.543 × 10⁻¹⁰ a⁻¹⁰ with a branching ratio of 0.1171 for $\lambda_{\varepsilon}/\lambda_{\beta t}$, but did not consider the uncertainties in these values. Despite its excellent potential age resolution the 40 Ar/ 39 Ar method as calibrated by these constants is limited in accuracy to no better than 2% arising mainly from the propagation of systematic errors in the 40 K decay constants and the 40 K $-{}^{40}$ Ar data for standards (Min et al., 2000).

Efforts to improve the accuracy of ⁴⁰Ar/³⁹Ar dating have focused on refining the ages of standards such as the widely used Fish Canyon sanidine (FCs), whose age has been addressed by the following approaches individually or in combination: (i) intercalibration with "first principles" standards (for which K-Ar data were available, e.g., Hurford and Hammerschmidt, 1985; Renne et al. 1998; Lanphere and Dalrymple, 2000; Spell and McDougall, 2003); (ii) comparison with ages based on orbital tuning (Renne et al., 1994; Hilgen et al., 1997; Kuiper et al., 2008); and (iii) comparison with results from coexisting mineral phases and/or different radioisotopic systems (e.g., Baksi et al., 1996; Villeneuve et al., 2000; Lanphere and Baadsgaard, 2001; Schmitz and Bowring, 2001; Bachman et al., 2007). The references cited above are intended to be indicative rather than exhaustive; for a more complete enumeration of data relevant to the age of FCs the reader is encouraged to consult those sources as well as the summary by Dazé et al. (2003).

The results of the various efforts to determine the age of FCs vary by several percent, far beyond what is expected from typical reported precision. Consensus is lacking as to the cause of this dispersion, but it certainly includes systematic errors and inconsistent interlaboratory calibrations as well as real geologic differences between the ages measured by different radioisotopic systems on different materials. Overall, the trend has been that older ages for FCs have been accepted in more recent times, culminating with the

unprecedentedly precise age of 28.201 ± 0.023 Ma reported by Kuiper et al. (2008) based on orbital tuning. The result of Kuiper et al. (2008) represented a significant advance in reconciling 40 Ar/ 39 Ar ages with those based on other chronometers including orbital tuning, and has achieved wide acceptance.

Despite the progress in determining the ages of standards presented by Kuiper et al. (2008), little progress has been made in refining the ⁴⁰K decay constants since the synthesis of Beckinsale and Gale (1969), updated by Steiger and Jäger (1977). As is discussed in the following, knowledge of the age of a standard in and of itself is only a partial solution to calibrating the ⁴⁰Ar/³⁹Ar system because miscalibrated decay constants accrue error as measured isotopic data are extrapolated to samples whose ages differ from that of the standard. Moreover, if we accept the K–Ar isotopic data for ⁴⁰Ar/³⁹Ar standards (i.e., as summarized by Jourdan and Renne, 2007), the only way to reconcile these data with the orbitally tuned result of Kuiper et al. (2008) is to adjust either or both of the principal ⁴⁰K decay constants.

Min et al. (2000) and Kwon et al. (2002) observed that one could improve the accuracy of determination of both λ_{tot} and the age of the FCs standard (or, in principle, any other standard) by comparing the ⁴⁰Ar/³⁹Ar ages of volcanic rocks with ages determined for the same samples by independent methods, in most cases by the U–Pb method. The U decay constants, particularly that for ²³⁸U, are significantly more precisely known than λ_{tot} or the age of FCs, so one can estimate these latter two quantities by finding values for them that best reconcile ⁴⁰Ar/³⁹Ar and ²³⁸U–²⁰⁶Pb* ages, as follows.

The ${}^{40}\text{Ar}*/{}^{39}\text{Ar}_{\text{K}}$ isotope ratio of a neutron-irradiated sample is related to these two parameters and to the true age of the sample by:

$$t_{\rm std} = \frac{1}{\lambda_{\rm tot}} \ln\left[\frac{(e^{\lambda_{\rm tot}\,\tau_i} - 1)}{R_i} + 1\right] \tag{1}$$

where λ_{tot} is the total ⁴⁰K decay constant, τ_i is an independently determined age of sample *i*, and R_i embodies the relationship between sample *i* and the standard, defined (Renne et al., 1998) by:

$$R_{i} \equiv \frac{({}^{40}\mathrm{Ar}^{*}/{}^{39}\mathrm{Ar}_{\mathrm{K}})_{i}}{({}^{40}\mathrm{Ar}^{*}/{}^{39}\mathrm{Ar}_{\mathrm{K}})_{\mathrm{std}}} \equiv \frac{(e^{\lambda_{\mathrm{tot}}\,t_{i}}-1)}{(e^{\lambda_{\mathrm{tot}}\,t_{\mathrm{std}}}-1)}$$
(2)

where ⁴⁰Ar^{*} denotes radiogenic ⁴⁰Ar and ³⁹Ar_K indicates ³⁹Ar produced from K by neutron irradiation. A U–Pb age (or other independent determination of the true age) and measurement of R_i for a particular sample thus define a relation between t_{std} and λ_{tot} . Kwon et al. (2002) considered five such pairs and applied a maximum likelihood method (discussed in more detail later) to find values of t_{std} and λ_{tot} that best fit these observations. Graphically, this can be thought of as determining the mutual intersection of a number of curves defined by Eq. (1) for (τ_i , R_i) pairs.

The approach of Kwon et al. (2002) showed the effectiveness of the basic idea of employing independently dated samples to calibrate the ${}^{40}\text{Ar}/{}^{39}\text{Ar}$ system, but did not incorporate constraints from activity measurements or isotopic data for the standard. In this work, these constraints

Independent parameter estimates used in the regression. See Sections 2.1-2.4 for discussion.				
Parameter	Value	Uncertainty (σ)	Relative Unc. (%)	Ref.
$\kappa_{\rm FCs}$	1.6408×10^{-3}	$0.0047 imes 10^{-3}$	0.289	This work

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$\kappa_{\rm FCs}$	1.6408×10^{-3}	$0.0047 imes 10^{-3}$	0.289	This work
$\lambda_{arepsilon}$	$0.580 imes 10^{-10} a^{-1}$	$0.007 imes 10^{-10} \mathrm{a}^{-1}$	1.207	Min et al. (2000)
λ_{β}	$4.884 imes 10^{-10} a^{-1}$	$0.049 \times 10^{-10} \mathrm{a}^{-1}$	1.003	Min et al. (2000)
$\lambda_{\rm tot}$	$5.5545 imes 10^{-10} \ a^{-1}$	$0.0109 imes 10^{-10} a^{-1}$	0.196	This work

are incorporated to determine t_{std} and both decay constants for ⁴⁰K. These parameters are related to (τ_i, R_i) pairs according to:

Table 1

$$\kappa_{\rm std} \equiv \left(\frac{{}^{40}{\rm Ar}^*}{{}^{40}{\rm K}}\right)_{\rm std} = \left(\frac{\lambda_{\varepsilon}}{\lambda_{\varepsilon} + \lambda_{\beta}}\right) \frac{\left[e^{(\lambda_{\varepsilon} + \lambda_{\beta})\tau_i} - 1\right]}{R_i} \tag{3}$$

in which Eq. (1) is transformed into a function of κ_{std} , the true (dimensionless) 40 Ar*/40 K ratio of the standard, rather than t_{std} . The (τ_i, R_i) pairs define relations between $\lambda_{\varepsilon}, \lambda_{\beta}$, and κ analogous to their relationship to t_{std} and λ_{tot} as employed by Kwon et al. (2002). We use a least-squares solution method similar to that of Kwon et al. (2002) to find values of λ_{ε} , λ_{β} , and κ that best fit (i) a set of such relations defined by seventeen (τ_i, R_i) pairs and (ii) independent determinations of λ_{ε} , λ_{β} , λ_{tot} , and κ_{std} . Graphically, this can be thought of as determining the mutual intersection of curved surfaces defined by Eq. (3) for the (τ_i, R_i) pairs, and planes defined by the independent measurements of λ_{ε} , λ_{β} , and κ . We show that the estimates of λ_{ε} , λ_{β} , and κ_{std} derived from this exercise (i) successfully reconcile U-Pb and ⁴⁰Ar/³⁹Ar ages for our entire data set, (ii) are significantly more precise than those determined from the existing independent measurements alone, and (iii) yield ⁴⁰Ar/³⁹Ar dates of unprecedented accuracy.

2. SUMMARY OF EXISTING DATA

A summary of existing constraints on λ_{ε} , λ_{β} , λ_{tot} , and the value of κ for the Fish Canyon sanidine (κ_{FCs}), are given in Table 1 and discussed below.

2.1. ⁴⁰Ar/⁴⁰K measurements of FCs

Three sources of data are available for the ⁴⁰Ar*/⁴⁰K ratio (κ_{FCs} value) of sanidine from the Fish Canyon Tuff. Steven et al. (1967) and Hurford and Hammerschmidt (1985) reported K-Ar data yielding mean ages of 27.9 ± 0.4 and 27.99 ± 0.27 Ma (respectively) based on measurements of $^{\rm 40}{\rm Ar}$ concentration by isotope dilution, K concentration by flame photometry, and the constants of Steiger and Jäger (1977). These K-Ar ages correspond to κ_{FCs} values of $(1.634 \pm 0.023) \times 10^{-3}$ and $(1.639 \pm 0.016) \times 10^{-3}$. Jourdan and Renne (2007) calibrated FCs against four primary standards (NL-25, Hb3gr, GHC-305, and GA1550) whose ⁴⁰K and ⁴⁰Ar were determined by isotope dilution using independently calibrated tracers. Equivalent K-Ar ages for FCs referred to these primary standards are 28.43 \pm 0.40, 28.03 ± 0.17 , 27.99 ± 0.13 , and 28.03 ± 0.16 Ma (respectively) yielding κ_{FCs} values of $(1.665 \pm 0.023) \times 10^{-3}$, $(1.641 \pm 0.010) \times 10^{-3}$, $(1.639 \pm 0.008) \times 10^{-3}$, and $(1.641 \pm 0.009) \times 10^{-3}$. The weighted mean $\kappa_{\rm FCs}$ value

from all six of these independent K-Ar or K-Ar equivalent measurements is $(1.6408 \pm 0.0047) \times 10^{-3}$, with an $MSWD^2$ of 0.25 and probability-of-fit of 0.94. Thus we conclude that the between-measurement scatter is consistent with stated uncertainties, and that the stated uncertainties may be overestimated.

2.2. β Activity data

Available data for the β activity of 40 K were evaluated by Min et al. (2000), who concluded that the previous compilation by Beckinsale and Gale (1969), which served as the basis for the decay constants recommended by Steiger and Jäger (1977), was flawed by arbitrary selection of activity data and, to a lesser extent, outdated physical constants. Using the more objective compilation of Endt and Van der Leun (1973), and updating the values of Avogadro's number and the atomic weight of K, Min et al. (2000) calculated a value of $(4.884 \pm 0.049) \times 10^{-10} a^{-1}$ for λ_{β} , which is used herein.

2.3. Electron capture activity data

Activity-based estimates of λ_{ε} are based on measurements of 1.46 MeV γ radiation. As with the β activity data, Min et al. (2000) used the compilation of Endt and Van der Leun (1973) and updated the values of physical constants to obtain a value of $\lambda_{\varepsilon} = (0.580 \pm 0.007) \times 10^{-10} \text{ a}^{-1}$. A hypothetical electron capture decay direct to ground state (without γ emission), accounting for 0.16% of total ^{40}K decay, was incorporated by Beckinsale and Gale (1969) and subsequently (implicitly) by Steiger and Jäger (1977). Lacking any evidence for the existence of this decay mode, we follow Min et al. (2000) in excluding it from consideration.

2.4. Liquid scintillation activity data

Two recent measurements of the total ⁴⁰K activity determined by liquid scintillation counting (LSC) techniques are available. Grau Malonda and Grau Carles (2002) measured solutions of KNO₃ with natural K, and reported a λ_{tot} value of $(5.554 \pm 0.009) \times 10^{-10} a^{-1}$. Kossert and Günther (2004) measured solutions of KNO₃ and KCl with natural K, and reported an identical value of $(5.554 \pm 0.013) \times$ $10^{-10} a^{-1}$. Combining all of the specific activities reported in these two studies with stoichiometric K concentrations yields a mean activity per atom of K of $(2.05445\pm0.00162)\times10^{-21}\,Bq/[K]$ (mean \pm standard error of the

² Mean square of weighted deviates (McIntyre et al., 1966).

mean), which then allows propagating the uncertainty in 40 K/K ((1.1672 ± 0.0021) × 10⁻⁴; Garner et al., 1975) as a systematic error. Thus, we obtain a value of $\lambda_{tot} = (5.5545 \pm 0.0109) \times 10^{-10} a^{-1}$ using the expression and constants (Avogadro's number, atomic weight of K, and seconds per tropical year) used by Min et al. (2000).

2.5. ²⁰⁶Pb/²³⁸U-⁴⁰Ar/³⁹Ar data pairs

The α activity measurements of ²³⁸U by Jaffey et al. (1971) provided an exceptionally precise value of λ (²³⁸U) = (1.55125 ± 0.00083) × 10⁻¹⁰ a⁻¹; consequently this nuclide has long been considered to present a logical basis for comparison or normalization of other decay constants (e.g., Wetherill et al., 1956; Steiger and Jäger, 1977; Patchett and Tatsumoto, 1980; Luck et al., 1980; Minster et al., 1982; Begemann et al., 2001; Schoene et al., 2006; Mattinson, in press). Note that, as pointed out by Begemann et al. (2001), many of the normalization studies (e.g., Patchett and Tatsumoto, 1980; Luck et al., 1980; Minster et al., 1982) used ²⁰⁷Pb/²⁰⁶Pb (rather than ²⁰⁶Pb/²³⁸U) ages as the normalization basis, which involves not only the ²³⁸U–²⁰⁷Pb decay, but also (and to a greater degree) that of ²³⁵U–²⁰⁷Pb. Thus ²⁰⁷Pb/²⁰⁶Pb ages are not only qualitatively different from ²⁰⁶Pb/²³⁸U ages, but also incur a substantial and generally overlooked penalty from the uncertainties in λ (²³⁸U) as well as λ (²³⁵U) (Ludwig, 2000). In this study only ²⁰⁶Pb/²³⁸U ages are used, as discussed further below.

The validity of our analysis rests in part on the validity of the value and uncertainty of the ²³⁸U decay constant reported by Jaffey et al. (1971). We share the conclusion of Schön et al. (2004) that further measurements of comparable precision are desirable in order to meaningfully evaluate the interlaboratory reproducibility of the Jaffey et al. (1971) results. Nonetheless we adopt the value and its uncertainty given by Jaffey et al. (1971). As we show subsequently, when we use these values we find that ²⁰⁶Pb/²³⁸U ages are consistent at their stated uncertainties with all measurements of λ_e , λ_β , λ_{tot} , and κ_{FCs} , which tends to validate the decay constant uncertainty of Jaffey et al. (1971). We then use pairs of ⁴⁰Ar/³⁹Ar and ²⁰⁶Pb/²³⁸U data from individual rocks, whose substitution into Eq. (3) define relations among κ_{FCs} , λ_e , and λ_β . The data pairs are shown in Table 2. The ⁴⁰Ar/³⁹Ar data are presented as *R* values relative to the Fish Canyon sanidine standard, as defined in Eq. (2) and based on data summarized in Table 3.

Surprisingly few pairs of reliable 206Pb/238U and $^{40}\text{Ar}/^{39}\text{Ar}$ data exist that permit valid comparison of these systems with temporal resolution approaching the precision of which the two systems are capable. The requirement of effectively simultaneous closure of both the U-Pb and K-Ar systems essentially limits the choices to precisely dated (precision <1%) volcanic products, whereas the vast majority of rocks dated by both methods are of plutonic or metamorphic origin (e.g., Schoene et al., 2006). Moreover, only volcanic materials free from reheating effects, as inferred from ⁴⁰Ar/³⁹Ar age spectrum and petrographic analysis, can be considered reliable. The volcanic origin criterion is particularly important for young (e.g., Phanerozoic) rocks wherein differential closure temperatures (Dodson, 1973) can produce relative age differences of several percent even at the geologically rapid cooling rates (>100 °C/Ma) typical of subvolcanic intrusions. In addition, the effects of preeruptive magma residence time of zircons must be considered even in volcanic rocks, so that pre-Mesozoic materials, whose age bias from pre-eruptive zircons should be less than 1_{00}° (e.g., Simon et al., 2008), are preferred.

Table 2

(τ_i, R_i) pairs used in the regression. See Appendix A1 for explanation and de
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Sample	τ^{a} (Ma)	σ_{τ} (Ma)	$\sigma_{ au}$ (%)	Ref.	R^{b}	σ_R	σ_R (%)	Ref.
79CE	1.919×10^{-3}			2	6.789×10^{-5}	$5.5 imes 10^{-7}$	0.81%	2
IP	17.42	0.03	0.17	3	0.60556	0.00249	0.41	3
P1T-2	125.65	0.17	0.14	1	4.5494	0.0061	0.13	4
PR94-7	133.06	0.26	0.20	1	4.8205	0.0153	0.32	1
NMB	201.27	0.13	0.07	5	7.4636	0.0072	0.10	6
MSG-09	242.14	0.45	0.19	7	9.0619	0.0122	0.13	7
SH-10	252.17	0.36	0.14	8	9.4983	0.0102	0.11	1,9
D3T	252.40	0.33	0.13	8	9.4918	0.0041	0.04	9
SH-09	252.60	0.33	0.06	8	9.5102	0.0059	0.06	1
PD97-2	253.02	0.50	0.20	1	9.4958	0.0122	0.13	10
SH-27	253.24	0.33	0.13	8	9.5295	0.0068	0.07	1
SH-16	253.68	0.30	0.12	8	9.5391	0.0085	0.09	1
SH-08	257.30	0.30	0.12	8	9.6987	0.0053	0.05	1
PaV	327.99	0.45	0.14	1	12.6039	0.0148	0.12	1
DK-LQ	454.59	0.56	0.12	1	18.0900	0.0523	0.29	11
F-239	1094.2	1.15	0.11	5	52.9011	0.1162	0.22	12
EGB-032	2067.5	1.4	0.07	5	135.6389	0.1821	0.13	1

Data sources: 1 – This study; 2 – Renne et al. (1998); 3 – Palfy et al. (2007); 4 – Chang et al. (2009); 5 – Schoene et al. (2006); 6 – Jourdan et al. (2009) 7 – Mundil et al. (2010); 8 – Mundil et al. (2004); 9 – Renne et al. (1995); 10 – Renne et al. (2001); 11 – Min et al. (2001); 12 – Min et al. (2000).

^a τ refers to independent age constraint, which is a calendar age for 79CE, and ²⁰⁶Pb/²³⁸U ages (uncorrected for residence time) for all others.

^b R refers to the relationship between ${}^{40}\text{Ar}*/{}^{39}\text{Ar}_{\text{K}}$ for the sample and that of the FCs standard, as defined by Eq. (2).

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Improved calibration of the K/Ar and Ar/Ar geochronometers

Table 3

Summary of new dat	a used to compi	le <i>R</i> -values, bas	ed on Ar isoto	pe data report	ed in EA-2.			
Sample	Lab. #	Funk	$\pm \sigma$	F _{std}	$\pm \sigma$	R	$\pm \sigma$	±σ (%)
PR94-7 Sanidine	30771-01 30771-02 30771-03 30771-04 30771-05 30771-06	11.55760 11.55990 11.56650 11.55650 11.55650 11.57490	$\begin{array}{c} 0.00190\\ 0.01200\\ 0.01050\\ 0.01100\\ 0.01130\\ 0.01210\\ \end{array}$					
FCs ^a	21181 21183 21185			2.40589 2.39648 2.39079	0.00259 0.00246 0.00317			
Summary		11.55820	0.00180	2.39772	0.00763	4.82050	0.01535	0.318
SH-16 Sanidine	33367	28.78470	0.01160					
FCs	33366			3.01755	0.00238			
Summary		28.78470	0.01160	3.01755	0.00238	9.53910	0.00846	0.089
SH-08 Sanidine	33010-11 33010-31 33010-32 33010-33	5.71136 5.71081 5.71196 5.71743	0.00385 0.00372 0.00369 0.00341					
FCs	33010			0.58906	0.00026			
Summary		5.71312	0.00183	0.58906	0.00026	9.69866	0.00532	0.055
SH-27 Sanidine	33016-11 33016-12 33016-13 33016-14 33016-15 33016-17 33016-19 33016-20 33016-22	5.56972 5.57344 5.57419 5.57406 5.56742 5.58538 5.56439 5.56634 5.56634	0.00733 0.00971 0.00883 0.00772 0.01134 0.00850 0.01201 0.01330 0.01059					
FCs	33017			0.58478	0.00026			
Summary		5.57269	0.00312	0.58478	0.00026	9.52952	0.00680	0.071
SH-10 Plagioclase	33018-21 33018-22 33018-23	5.57037 5.58223 5.58423	0.02255 0.01006 0.01044					
FCs	33019			0.58743	0.00050			
Summary		5.58199	0.00690	0.58743	0.00050	9.50237	0.01428	
Previous Data (^b)						9.49405	0.01464	
Global Summary						9.49831	0.01022	0.108
SH-09 Sanidine	33020-11 33020-12	5.59190 5.59870	0.00337 0.00398					
FCs	33021			0.58813	0.00027			
Summary		5.59318	0.00241	0.58813	0.00027	9.51018	0.00594	0.062
PaV Hornblende	32451-01 32541-02 32541-03 32451-10 32451-13	37.35305 37.30356 37.19669 37.05151 37.00940	0.04465 0.05521 0.26240 0.14302 0.15900					
FCs	32452			2.95963	0.00230			
Summary		37.30280	0.03270	2.95963	0.00230	12.60389	0.01475 (continued on	0.117 next page)

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Table 2 (continued)

Table 5 (commuted)							
Sample	Lab. #	$F_{\rm unk}$	$\pm \sigma$	F _{std}	$\pm \sigma$	R	$\pm \sigma$	±σ (%)
EGB-032	33000-01	79.72397	0.54839					
Hornblende	33000-03	79.64110	0.41550					
	33000-05	79.92520	0.21257					
	33000-06	79.79686	0.26976					
	33000-07	79.99423	0.21664					
	33000-08	79.85893	0.32670					
	33000-11	79.57674	0.21405					
	33000-14	79.90703	0.26277					
	33000-15	80.26145	0.20002					
	33000-18	79.57250	0.46829					
	33000-20	80.21144	0.40518					
FCs	33001			0.58912	0.00051			
Summary		70 00788	0.08231	0.58012	0.00051	135 63803	0 18206	0.134

Notes. Samples are described in Appendix A2. *F* values refer to ${}^{40}\text{Ar}^*/{}^{39}\text{Ar}_{K}$; *R* values as defined in text. Summary values are weighted (by inverse variance) means. FCs data are from grains in the same irradiation disk pit as the sample in all but the case of PR94-7 (^a), where FCs from three pits bracketing the sample were averaged and the mean and standard deviation are used. For sample SH-10, the Previous Data (^b) are from Renne et al. (1995).

To be included in our analysis, ²⁰⁶Pb/²³⁸U ages must be concordant with ²⁰⁷Pb/²³⁵U ages at the 95%-confidence level using the concordia-age algorithms of Ludwig (1998), specifically including decay constant uncertainties. Normalization of ⁴⁰Ar/³⁹Ar ages in this study is restricted to the ²³⁸U-²⁰⁶Pb decay system, as recommended by Steiger and Jäger (1977), eliminating the effects of ²³⁵U decay-constant uncertainties (Mattinson, 1987; Ludwig, 1998, 2000) and/or open-system behavior on the uncertainties of ²⁰⁷Pb/²⁰⁶Pb and concordia-intercept ages. ²⁰⁶Pb/²³⁸U data from multiple single-zircon analyses must be statistically coherent at the 95%-confidence level to avoid (as much as possible) averaging complexities arising from Pb-loss and/or subtle inheritance of pre-eruptive crystals. Pretreatment of most zircon samples with the anneal-leach technique pioneered by Mattinson (2005) produced results consistent with closed isotopic systems, eliminating a need for arbitrary selection from complex data sets.

To be used for our analysis, 40 Ar/ 39 Ar data must exhibit excellent reproducibility and demonstrable freedom from the effects of secondary alteration or 39 Ar recoil. Strict attention to determining the neutron fluence, and accounting for possible fluence gradients during sample irradiation, is required. Non-atmospheric initial 40 Ar must be absent or precisely characterized by isochron analysis. Data from pyroclastic rocks must be obtained from single crystals to avoid xenocrystic contamination. Age spectra obtained by incremental heating must define a plateau (*sensu* Fleck et al., 1977) containing >90% of the 39 Ar released and lacking statistically significant slope.

U–Pb and ⁴⁰Ar/³⁹Ar analyses were performed at the Berkeley Geochronology Center using facilities and procedures described previously (Min et al., 2000; Mundil et al., 2001, 2004) except as noted, with results summarized in Electronic Annexes EA-1 and EA-2 respectively. Descriptions of samples lacking previously published data are given in the Appendix. Some additional data from the literature which meet our criteria are also included. The volcanic rocks span an age interval between ca. 17 and 2068 Ma. ²⁰⁶Pb/²³⁸U ages include uncertainty in the ²³⁸U decay constant, tracer calibration, and correction for initial ²³⁰Th. ⁴⁰Ar/³⁹Ar results are represented by *R*-values as defined in Eq. (2) relative to Fish Canyon sanidine. The results used to compute weighted (by inverse variance) mean *R*-values from ⁴⁰Ar/³⁹Ar data presented in Electronic Annex EA-2 are summarized in Table 3. ⁴⁰Ar/³⁹Ar ages corresponding to these *R*-values, using the decay constants of Steiger and Jäger (1977) and the standard calibration of Renne et al. (1998), hereinafter referred to as "conventional" ⁴⁰Ar/³⁹Ar ages, are compared with the ²⁰⁶Pb/²³⁸U ages in Fig. 1.

The data pairs of ²⁰⁶Pb/²³⁸U and conventional 40 Ar/ 39 Ar ages presented here exhibit a systematic bias such that the ${}^{206}\text{Pb}/{}^{238}\text{U}$ ages are older than the ${}^{40}\text{Ar}/{}^{39}\text{Ar}$ ages. The bias ranges from >1.5% for young rocks to ca. 0.5% for rocks as old as 2 Ga. A bias of this magnitude resulting from miscalibrated artificial tracers used in isotope dilution U/Pb analyses is not expected as recent cross-calibrations of the U/Pb ages presented here are comparable at the 0.15% level with ages from other state-of-the-art laboratories (Black et al., 2004). A potential variation of λ_{ε} due to changing physical and chemical environments has been shown to be even less significant (Bukowinski, 1979; Norman et al., 2001). Instead, the magnitude and observed pattern of data pairs (bias decreasing as a function of age) can most easily be explained by two independent contributing factors: pre-eruptive residence time and mis-estimation of the ⁴⁰K decay constants.

Pre-eruptive magma residence times for zircons, which have been observed to be as large as ca. 300 ka, can have a significant impact on relatively young rocks (such as sample IP of this study), but their importance should diminish linearly with age. In other words, if residence time was the sole reason for the bias in all samples depicted in Fig. 1(a), the observed bias of 400 ka to 13 Ma could only be explained by a correlation between age and residence time. This is implausible not only because there is no physical mechanism to produce such a correlation, but also because



Fig. 1. Bias between ${}^{238}U/{}^{206}Pb$ zircon ages and "conventional" calibration of the ${}^{40}Ar/{}^{39}Ar$ system for data pairs listed in Table 2. Upper panel (a) shows ${}^{238}U/{}^{206}Pb$ zircon ages uncorrected for residence time effects. Lower panel (b) shows the effect of correcting for residence time of 90 \pm 77 ka (see text for discussion). For comparison are shown curves depicting the effects of: (i) 100–1500 ka zircon residence times (solid gray); (ii) decreasing the electron capture decay constant for ${}^{40}K$ by 0.5%, 1.0%, and 1.5% (dashed black); and (iii) similarly decreasing the beta decay constant of ${}^{40}K$ (solid black).

zircon residence times greater than 1 Ma would be required for all of the >100 Ma samples – much longer than any case documented to date. Inherited zircons or antecrysts (e.g., Bacon and Lowenstern, 2005) exhibiting such pre-eruptive age are not uncommon, but in cases considered here are ruled out by the statistical coherence of the individual zircon $^{206}Pb/^{238}U$ ages.

In Fig. 1(b) we recalculated the ${}^{206}\text{Pb}/{}^{238}\text{U}$ ages assuming that they all include a pre-eruptive residence time of 90 ± 77 ka derived from the compilation of Simon et al.

(2008). Although further studies are desirable to clarify the scope of residence time manifest in zircon U–Pb ages, the correction we use is an objective measure given the current state of knowledge. The recalculation consists of simply subtracting 90 ± 77 ka from each ²⁰⁶Pb/²³⁸U age and propagating the additional uncertainty such that if the uncorrected ²⁰⁶Pb/²³⁸U age uncertainty is σ_{τ} , the uncertainty on the corrected ²⁰⁶Pb/²³⁸U age σ'_{τ} becomes $\sqrt{\sigma_{\tau}^2 + (77 \text{ ka})^2}$. This adjustment slightly reduces ²⁰⁶Pb/²³⁸U ages, and slightly

increases their uncertainties. As discussed later, we adopt this correction to the ²⁰⁶Pb/²³⁸U ages and their uncertainties in the regression analysis. After accounting for residence time bias in this way, there is still a tendency for the bias to decrease as a function of age although all the individual results are indistinguishable at 95% confidence from a 1% bias. This tendency is easily explained in sense and magnitude by miscalibration of λ_{ε} , whose relative uncertainty from disintegration counting is the largest among the three variables considered here.

Whereas the ca. -1% bias implying miscalibration of λ_{e} is well constrained for Phanerozoic samples, a decrease of the bias to ca. -0.5% (possibly even lower) for older ages (as displayed by the data pairs from F-239 and EGB-032 at 1.09 and 2.06 Ga, respectively), is less confidently attributed to this source. Rather, miscalibration of λ_{β} in the opposite sense is suggested, consistent with the previously discussed LSC data which indicate a $\sim 0.22\%$ larger value ($5.555 \times 10^{-10} a^{-1}$) of λ_{tot} than that ($5.543 \times 10^{-10} a^{-1}$) recommended by Steiger and Jäger (1977). This hypothesis is addressed more thoroughly by the regression analysis described later.

Several data sets are excluded from our analysis that might seem appropriate to include. Specifically, U-Pb zircon data from the Fish Canyon Tuff (FCT) itself could be highly relevant in principle to our analysis. Unfortunately, such data from the FCT (Schmitz and Bowring, 2001; Bachman et al., 2007) and two additional eruptive units from the same magma system (Bachman et al., 2007) reveal excess dispersion (i.e., having MSWD values ranging from 6 to 18) in 206 Pb/ 238 U, and thereby fail the equivalence (sensu Ludwig, 1998) criterion employed to exclude data affected by Pb-loss and/or subtle inheritance. We note here that although Schmitz and Bowring (2001) reported a seemingly reasonable MSWD of 0.97 for their FCT zircon data, Bachman et al. (2007) showed that this was an artifact of a calculation error, and that the true MSWD is ~ 10 . The excess scatter manifest in U-Pb data for all three units from this magma system probably arises – at least in part – from pre-eruptive residence time, but the observed non-analytical scatter in the individual dates precludes any confidence that Pb-loss has been completely mitigated, or that the effects of subtle inheritance and/or antecrysts are absent. A similar caveat applies to ${}^{206}\text{Pb}/{}^{238}\text{U} - {}^{40}\text{Ar}/{}^{39}\text{Ar}$ data pairs reported for tuffs in the Green River Formation (Smith et al., 2010), wherein excess scatter in ²⁰⁶Pb/²³⁸U ages violates our equivalence criterion and subjective data culling is required to interpret an age. Therefore these data, despite their excellent analytical quality, are excluded from our analysis although we note that excess scatter of similar magnitude could be present (though unrecognizable) in older samples including those used in this study.

2.6. Other (τ_i, R_i) pairs

In addition to samples with ${}^{40}\text{Ar}/{}^{39}\text{Ar}$ data whose age is independently constrained by ${}^{206}\text{Pb}/{}^{238}\text{U}$ dates, there are several additional data pairs possible in which the independent constraints are provided by other means. Renne and Min (1998) proposed such an approach using ${}^{40}\text{Ar}/{}^{39}\text{Ar}$ data for sanidine from the historically dated Vesuvius eruption of 79 CE (Common Era). Although Renne and Min (1998) used a different standard (ACs) as a neutron fluence monitor, their results can be related to FCs using the definitions and data given by Renne et al. (1998), such that $R_{FCs}^{AD79} \equiv R_{ACs}^{AC3} \cdot R_{FCs}^{ACs} = (6.789 \pm 0.055) \times 10^{-5}$, as used by Kwon et al. (2002). The relevant age for comparison (in CE 1998, when the measurements were made or normalized to) is 1.919 ka, and we include this data pair in our analysis despite the relative imprecision ($\pm 0.8\%$) of the *R*-value.

Another set of (τ_i, R_i) pairs that might be included are those from Kuiper et al. (2008) that consist of ⁴⁰Ar/³⁹Ar data relative to FCs for Miocene tuffs in the Melilla Basin of Morocco whose absolute ages can be inferred from astronomical tuning, that is, correlating stratigraphic climate proxy records with calculated orbitally modulated terrestrial insolation time series. These data would complement the ²⁰⁶Pb/²³⁸U-constrained data pairs well because they are relatively young, but not affected by pre-eruptive residence time. However, because ages derived from astronomical tuning are based on one-to-one correlation of stratigraphic units with insolation cycles rather than on analyses with well-defined uncertainty distributions, it is difficult to estimate their uncertainty in a way that is compatible with measurement uncertainties on our other input data. Thus, we exclude these results from the present analysis and instead, as discussed later, use them as a basis for independent comparison.

3. DATA ANALYSIS

We seek to find values for the parameters λ_{ε} , λ_{β} , and κ_{FCs} that best fit: (i) the relations described by Eq. (3) and the (τ_i , R_i) pairs described in the Introduction (Section 1) and Section 2.5, as well as (ii) the independent measurements of these parameters described in Sections 2.1–2.4. We adapt the nonlinear regression model of Kwon et al. (2002) for this purpose. First, we rewrite Eq. (3) to predict a value for the true age of a sample τ_i^* given estimates of R_i^* , λ_{ε}^* , λ_{β}^* , and κ_{FCs}^* . The superscripted star denotes estimates for the various quantities that are generated by the solution method, as opposed to the independently measured values of these quantities that comprise the data set we seek to fit. This yields:

$$\tau_i^* = \left(\frac{1}{\lambda_{\varepsilon}^* + \lambda_{\beta}^*}\right) \ln\left[\left(\frac{\lambda_{\varepsilon}^* + \lambda_{\beta}^*}{\lambda_{\varepsilon}^*}\right) \kappa_{\text{FCs}}^* R_i^* + 1\right]$$
(4)

Following Kwon et al. (2002), the best estimates of λ_{ε}^* , λ_{β}^* , and κ_{FCs}^* are found by minimizing the sum of squares *S*:

$$S(X) = \sum_{i=1}^{n} \left[\frac{1}{2} \left(\frac{\tau_i^* - \tau_i}{\sigma_{\tau,i}} \right)^2 + \frac{1}{2} \left(\frac{R_i^* - R_i}{\sigma_{R,i}} \right)^2 \right]$$
(5)

over $X = (\lambda_{\varepsilon}^*, \lambda_{\beta}^*, \kappa^*, R_1^*, R_2^*, \dots, R_n^*)$. Here $\sigma_{\tau,i}$ and $\sigma_{R,i}$ are measurement uncertainties in τ_i and R_i , respectively (Table 2). Note that this procedure also yields optimized estimates of the *R* values of all the samples in our data set, though these are nuisance parameters for the present purpose. Minimizing S(X) yields parameter values that best

fit the (τ_i, R_i) pairs, but our overall goal in this work is to incorporate all pertinent information, including both the (τ_i, R_i) pairs and the independent measurements of κ_{FCs} and the decay constants, into our estimates of these quantities. Thus, we find values of λ_{ε}^* , λ_{β}^* , and κ_{FCs}^* that best fit both the (τ_i, R_i) pairs and the independent measurements λ_{ε} , λ_{β} , λ_{tot} , and κ_{FCs} by minimizing

$$S_{1}(X) = S(X) + \left(\frac{\lambda_{\varepsilon}^{*} - \lambda_{\varepsilon}}{\sigma_{\lambda\varepsilon}}\right)^{2} + \left(\frac{\lambda_{\beta}^{*} - \lambda_{\beta}}{\sigma_{\lambda\beta}}\right)^{2} + \left(\frac{\kappa_{\text{FCs}}^{*} - \kappa_{\text{FCs}}}{\sigma_{\text{K}_{\text{FCs}}}}\right)^{2} + \left(\frac{\lambda_{\delta}^{*} - \lambda_{\text{tot}}}{\sigma_{\lambda_{\text{tot}}}}\right)^{2}$$
(6)

over $X = (\lambda_{\varepsilon}^{*}, \lambda_{\beta}^{*}, \kappa_{FCs}^{*}, R_{1}^{*}, R_{2}^{*}, \dots, R_{n}^{*})$. Here $\lambda_{\varepsilon}, \lambda_{\beta}, \lambda_{tot}$, and κ_{FCs} (lacking superscripted asterisks) are the independently measured values of these quantities described above in Section 2, and $\sigma_{\lambda\varepsilon}, \sigma_{\lambda\beta}, \sigma_{\lambda tot}$, and $\sigma_{K_{FCs}}$ are their respective measurement uncertainties (Table 1). An important aspect of this approach is that the independent estimate of each of these parameters is given the same weight as any one of the (τ_i, R_i) pairs.

As previously discussed in Section 2, we adjusted all values of τ_i in Table 2 that are derived from ²⁰⁶Pb/²³⁸U measurements by subtracting a pre-eruptive residence time of 90 ± 77 ka. Although this correction acts to slightly improve the optimal value of S_1 , we applied it not for this reason but because of the first-principles observation that preeruptive residence time is common and, therefore, that a correction is required. We then actually carried out the minimization using the MATLAB implementation of the Nelder-Mead simplex search method (Lagarias et al., 1998), and estimated uncertainties in the best-fitting parameter values by a 4000-trial Monte Carlo simulation. With one exception, all measurement uncertainties in the input parameters were assumed to be Gaussian and independent. The exception is that the uncertainties in the ${}^{206}Pb/{}^{238}U$ ages are not independent, because of their common dependence on the ²³⁸U decay constant. We incorporated this effect into the Monte Carlo simulation by: (i) transforming ²⁰⁶Pb/²³⁸U ages and their uncertainties to ratios of radiogenic ²⁰⁶Pb to ²³⁸U (²⁰⁶Pb*/²³⁸U) and corresponding uncertainties using the Jaffey et al. (1971) value and uncertainty for λ (²³⁸U); (ii) for each trial, randomly generating values of ²⁰⁶Pb*/²³⁸U for all samples as well as a single common value of λ (²³⁸U); and (iii) using these to compute trial values of τ_i . This procedure produces a random sample of τ_i values that display the covariance resulting from their common dependence on the ²³⁸U decay constant.

4. OPTIMIZATION RESULTS

Tables 4a–c show the values of λ_{ε}^* , λ_{β}^* , κ_{FCs}^* , and R_1^* , R_2^* , ..., R_n^* that minimize S_1 , as well as the implied values of τ_i^* . These best-fit values yield $S_1 = 8.8$. Table 4, Figs. 2 and 3 show the results of the Monte Carlo error analysis. Fig. 4 shows the distribution of the normalized residuals with respect to the input values of τ_i , R_i , λ_{ε} , λ_{β} and κ_{FCs} ; the distribution is indistinguishable from normal. We did not observe any significant relationship between the



Fig. 2. Optimal values of the parameters κ_{FCs} , λ_e , and λ_β , and results of the Monte Carlo error analysis compared with constraints used as input to the regression analysis. Each panel has two of these parameters as its axes, and the constraints are evaluated at the optimal value, given in Table 4a, of the third parameter. For example, in the panel with axes κ_{FCs} and λ_e , the constraints are computed with λ_β fixed at $4.9737 \times 10^{-10} a^{-1}$. The open circles show the optimal values of the parameters. The black dots show projections normal to the plane of each panel of the results of the Monte Carlo error analysis. The blue lines and blue shaded regions represent planes and corresponding 68% confidence regions defined by the independent parameter measurements in Table 1. The red lines and red shaded regions represent surfaces and corresponding 68% confidence regions defined by Eq. (3) evaluated for each of the (τ_i , R_i) pairs in Table 3.



Fig. 3. Marginalized probability distributions for the parameters κ_{FCs} , λ_e , and λ_β derived from the Monte Carlo error analysis. The dark vertical line shows the optimal value of each parameter, and the grey band shows the standard deviation of each parameter computed from the Monte Carlo results.



Fig. 4. Frequency distribution of normalized residuals for the bestfitting values of λ_{e}^{*} , λ_{β}^{*} , κ_{FCs}^{*} , R_{1}^{*} , R_{2}^{*} , ..., R_{n}^{*} , and the implied values of τ_{i}^{*} , and the corresponding measurements and their uncertainties.

residuals with respect to R_i and τ_i and the values of R_i and τ_i themselves. One can evaluate the quality of the fit by computing, for example, the reduced chi-squared (χ_R^2) statistic of the best-fitting values of λ_{ε}^* , λ_{β}^* , κ_{FCs}^* , R_1^* , R_2^* , ..., R_n^* , and the implied values of τ_i^* with respect to the corresponding measurements and their uncertainties: here $\chi^2_R = 0.75$ for 18 degrees of freedom. This is close to the expected value of 1, with a probability-of-fit of 0.76, indicating that the scatter of the measurements with respect to the best-fitting parameter values is consistent with measurement uncertainties. We conclude that: (i) these best-fitting values for λ_{ε}^* , λ_{β}^* , and κ_{FCs}^* successfully reconcile all paired ${}^{40}\text{Ar}/{}^{39}\text{Ar}$ and $^{206}\mbox{Pb}/^{238}\mbox{U}$ ages as well as the independent measurements of the parameters and (ii) all relevant measurements incorporated in the regression analysis are consistent at the level of their stated uncertainties.

5. APPLICATION OF THE RESULTS

Comparison of Tables 1 and 4a shows that our estimates of λ_{ε}^* and λ_{β}^* are significantly more precise than previously determined values, while the precision of our estimate of κ_{FCs}^* remains essentially unchanged.

5.1. Age uncertainties

We now draw attention to two additional aspects of the results that are critically important in using these parameter estimates to compute 40 Ar/ 39 Ar ages and their uncertainties. First, it is important to make clear that because we are estimating all three parameters needed for computing 40 Ar/ 39 Ar ages simultaneously, these parameter estimates are not independent. In fact, as shown by the Monte Carlo results in Fig. 2, the uncertainties on these parameters are strongly correlated. This contrasts with the previous measurements shown in Table 1, which were determined in separate experiments with independent uncertainties. One must take the covariance into account in error propagation when calculating 40 Ar/ 39 Ar ages using our estimates of λ_{e}^{*} , λ_{β}^{*} , and κ_{FCs}^{*} . Using linear error propagation, the uncertainty in an 40 Ar/ 39 Ar age calculated using Eq. (4) is:

$$\begin{aligned} \sigma_{t}^{2} &= \left(\frac{\partial t}{\partial \lambda_{\varepsilon}} \sigma_{\lambda_{\varepsilon}}\right)^{2} + \left(\frac{\partial t}{\partial \lambda_{\beta}} \sigma_{\lambda_{\beta}}\right)^{2} + \left(\frac{\partial t}{\partial \kappa_{\rm FCs}} \sigma_{\kappa_{\rm FCs}}\right)^{2} \\ &+ \left(\frac{\partial t}{\partial R} \sigma_{R}\right)^{2} + 2\frac{\partial t}{\partial \lambda_{\varepsilon}} \frac{\partial t}{\partial \kappa_{\rm FCs}} \operatorname{cov}(\lambda_{\varepsilon}, \kappa_{\rm FCs}) + 2\frac{\partial t}{\partial \lambda_{\beta}} \\ &\times \frac{\partial t}{\partial \kappa_{\rm FCs}} \operatorname{cov}(\lambda_{\beta}, \kappa_{\rm FCs}) + \frac{\partial t}{\partial \lambda_{\varepsilon}} \frac{\partial t}{\partial \lambda_{\beta}} \operatorname{cov}(\lambda_{\varepsilon}, \lambda_{\beta}) \end{aligned} \tag{7}$$

Table 4a

Optimization results: optimal values of parameters needed to compute a ${}^{40}\text{Ar}/{}^{39}\text{Ar}$ age $(K, \lambda_{c}\lambda_{\beta})$ as well as the value of λ_{tot} implied by these results. The uncertainties are derived from the Monte Carlo error analysis.

Variable	Value	$\pm \sigma$	±%
KEC.	1.6418×10^{-3}	0.0045×10^{-3}	0.27
λ_{e}	$0.5755 \times 10^{-10} a^{-1}$	$0.0016 \times 10^{-10} a^{-1}$	0.28
λ_{B}	$4.9737 \times 10^{-10} \ a^{-1}$	$0.0093 \times 10^{-10} \ a^{-1}$	0.19
$\dot{\lambda}_{tot}$	$5.5492 \times 10^{-10} \ a^{-1}$	$0.0093 imes 10^{-10} a^{-1}$	0.17

Table 4b Optimization results: covariances and error correlations between parameters needed to compute a ⁴⁰Ar/³⁹Ar age (*K*, $\lambda_e \lambda_\beta$), derived from the Monte Carlo error analysis.

Variables (x,y)	Covariance $cov(x, y)$	Error correl. $\rho_{x,y}$		
$ \begin{array}{c} (\kappa_{\rm FCs},\lambda_{\varepsilon}) \\ (\kappa_{\rm FCs},\lambda_{\beta}) \\ (\lambda_{\varepsilon},\lambda_{\beta}) \end{array} $	$\begin{array}{c} 7.1889 \times 10^{-19} \\ -7.1390 \times 10^{-19} \\ -3.4497 \times 10^{-26} \end{array}$	0.9985 -0.1706 -0.2318		

This includes covariance terms for λ_{ε} , λ_{β} , and κ_{FCs} . A measurement of R is independent of these parameters, so the covariance terms involving R are zero and are omitted.Solutions of the partial derivatives in Eq. (7) are given in Appendix A.2 and the covariances are as given in Table 4b along with the corresponding error correlations defined by $\rho_{x,y} = \operatorname{cov}(x, y)/(\sigma_x \sigma_y)$. The covariance terms in this expression have a significant effect on the uncertainty of an ${}^{40}\text{Ar}/{}^{39}\text{Ar}$ age computed from the data in Table 4a. For example, a hypothetical sample with $R = 50 \pm 0.25$ (a 0.5% uncertainty) has an age of 1050.7 Ma. The uncertainty in this age calculated (incorrectly) ignoring the covariance terms in Eq. (7) would be ± 5.1 Ma; the correct uncertainty that takes account of the covariance terms is ± 4.0 Ma. A graphical explanation of this effect is as follows. Finding the values of λ_{ε}^* , λ_{β}^* , and κ_{FCs}^* that minimize Eq. (5) is in effect finding values of these parameters that best reconcile ⁴⁰Ar/³⁹Ar and U/Pb ages. Graphically, this can be thought of as seeking points that are near the mutual intersection of the surfaces defined by Eq. (3) applied to the various (τ_i, R_i) pairs. As is evident from Fig. 2, these intersections are

nearly colinear in the region of interest. Thus, parameter values can be relatively widely scattered with respect to any one of the individual parameters, but if they remain close to this region of intersection they will still precisely reconcile 40 Ar/ 39 Ar and U/Pb ages.

A second important issue for applying our results is that the uncertainties and covariances reported in Tables 4a and b are derived via Monte Carlo simulation, so are not arbitrarily precise. In addition, Eq. (7) above is an approximation derived by linearizing Eq. (4). For these reasons, Eq. (7) with the uncertainties and covariances in Tables 4a and b may yield age uncertainties that are inaccurate when uncertainties in R are less than approximately 1 per mil. Thus, one should use a Monte Carlo method rather than linear error propagation to compute age uncertainties when the uncertainties in R are small. We supply in EA-3 an Excel spreadsheet containing (i) 4000 (λ_{ε} , λ_{β} , and κ_{FCs}) triplets, that are derived from our Monte Carlo uncertainty analysis for these parameters and display the variances and covariances shown in Tables 4a and b, and (ii) a spreadsheet program for generating a normal distribution of *R*-values given an arbitrary input mean and standard deviation. This spreadsheet can be used straightforwardly for a Monte Carlo analysis of uncertainty in age calculations. Note that for computing ages themselves, the exact solution presented by Eq. (4) should be used; the age returned by the Monte Carlo approach will always be an approximation.

Relative age errors estimated by Eq. (7) are shown in Fig. 5 for values of R with relative precision between 0% and 1.0%. A noteworthy consequence of our results, illustrated in Fig. 5, is that the relative uncertainty of 40 Ar/ 39 Ar

Fig. 5. Summary of the accuracy of 40 Ar/ 39 Ar ages calculated with the results of this work, compared with the 238 U/ 206 Pb system. Relative age errors are shown for different values of relative isotope measurement precision (σ_R/R) whose values correspond to the intercept on the ordinate (relative age error) axis; lowest (darker) curves correspond to infinitely precise *R*-values (zero uncertainty). For the 40 Ar/ 39 Ar system, *R* is as defined by Eq. (2) as the 40 Ar*/ 39 Ar_K of a sample relative to that of the FCs standard. For the U/Pb system, *R* is defined as 206 Pb*/ 238 U.



ages is dramatically improved with increasing age (*R*-value) for relative σ_R greater than about 0.2%. While contraction in relative age error with increasing age is a consequence of the logarithmic relation between age and R, and is true for all radioisotopic age equations, the relationship shown in Fig. 5 is especially noteworthy in comparison with the 206 Pb/ 238 U system. Here, the *R*-value for 206 Pb/ 238 U is taken to be the ²⁰⁶Pb*/²³⁸U ratio, and its uncertainty includes contributions from blank and mass fractionation corrections, spike calibration, etc. For ⁴⁰Ar/³⁹Ar ages, the *R*-values in Fig. 5 are as defined by Eq. (2) and also include corrections for blank and mass fractionation as well as decay and nucleogenic interferences. Thus the R-values and their uncertainties used for comparison of the two systems are not strictly equivalent, but they are comparable measures of the relevant mass spectrometric data and their accuracy.

An obvious implication of Fig. 5 is that with comparable measurement precision, the ${}^{40}\text{Ar}/{}^{39}\text{Ar}$ system can, with existing techniques, produce ages with accuracy comparable to or even exceeding the ${}^{206}\text{Pb}/{}^{238}\text{U}$ system. In effect this potential can now be realized because we have calibrated the ${}^{40}\text{Ar}/{}^{39}\text{Ar}$ system to both the ${}^{238}\text{U}$ and ${}^{40}\text{K}$ decay constants. This does not suggest that the ${}^{40}\text{Ar}/{}^{39}\text{Ar}$ method is necessarily preferable to the U/Pb method in cases where both are applicable, such as cases in which daughter-element mobility in the ${}^{40}\text{Ar}/{}^{39}\text{Ar}$ system is an issue, but it places the ${}^{40}\text{Ar}/{}^{39}\text{Ar}$ method on firm footing to validate even exceptionally precise U/Pb ages in cases where both are expected to record simultaneous closure to daughter nuclide mobility. Some examples follow.

5.2. Comparison with astronomical calibration

The previously-mentioned recent astronomical calibration of the age of the FCs standard (Kuiper et al., 2008) provided a significant breakthrough in reducing systematic errors and reconciling ⁴⁰Ar/³⁹Ar ages with ages determined by U/Pb and astronomical tuning. Application of the results of Kuiper et al. (2008) is based on Eq. (1) and the relationship between an unknown and the suronomicallydated horizon given by $R_{Ast}^{un} \equiv R_{FCs}^{un} \cdot R_{Ast}^{FCs} = R_{FCs}^{unk} \cdot (4.3644 \pm 0.0009)$. In this case the independent "true" age τ (6.500 \pm 0.005 Ma) is determined independently by astronomical tuning as discussed in Section 2.6.

In comparing our results to the calibration of Kuiper et al. (2008), which was focused on the age of FCs, we first address the question of how ${}^{40}\text{Ar}/{}^{39}\text{Ar}$ ages based on this result compare to ages and age uncertainties computed using our calibration. We note that Kuiper et al. (2008) did not explicitly recommend a value for λ_{tot} , but their derivation of an age for FCs used the sum of the partial decay constants listed in Table 1, $(5.463 \pm 0.107) \times 10^{-10} \text{ a}^{-1}$, as described above. The λ_{tot} dependence of Eq. (1) introduces large propagated age errors for astronomically calibrated ages very different from τ_{Ast} . These age errors are significantly larger than those computed using our new calibration as shown in Fig. 6.

Ages computed from the astronomical calibration as described above are compared with those based on our new calibration in Fig. 7. The values of λ_{ε} , λ_{β} , and κ_{FCs} determined herein yield older ages (by 0.38%) at zero age, and younger ages (by 0.86%) at 4560 Ma than those computed using the astronomical calibration. However, due mainly to the large uncertainties in the latter, the differences between ages based on the two calibrations are not significant at 95% confidence except for ages <100 Ma with precisely determined *R*-values ($\sigma_R/R < 0.1\%$).

Second, we ask whether the (τ, R) pair itself determined by the combination of astronomical tuning and ${}^{40}\text{Ar}/{}^{39}\text{Ar}$ analysis relative to FCs is consistent with our independent calibration. To do this, we used the relationship $R_{\text{FCs}}^{\text{Ast}} = (R_{\text{Ast}}^{\text{FCs}})^{-1} = 0.229127 \pm 0.000047$ to compute an age based our values of λ_e , λ_β , and κ_{FCs} . The resulting age is



Fig. 6. Age uncertainties as a function of age for calibration of the 40 Ar/ 39 Ar system presented herein versus that of Kuiper et al. (2008) based on astronomical tuning for the age of FCs and the 40 K decay constant of Min et al. (2000). Contours show results for various relative uncertainties in *R*.



Fig. 7. Difference between ages determined from data presented herein versus those based on the astronomical age of Kuiper et al. (2008) for FCs.

 6.525 ± 0.005 Ma (error computed using Monte Carlo simulation), which can be compared with the astronomically tuned age of 6.500 ± 0.005 Ma reported by Kuiper et al. (2008). The difference between these ages, 25 ± 7 ka, is significant at 95% confidence. Though this age difference is indistinguishable from the duration of a single orbital precession cycle, this may be serendipitous as the astronomical tuning of Kuiper et al. (2008) is based on orbital eccentricity and obliquity as well as precession, whose combined effect on insolation is the tuning basis. The cause of the discrepancy should be addressed by future work.

5.3. Implications for selected ⁴⁰Ar/³⁹Ar ages

Table 5

Several implications of the revised constants proposed here bear discussion. We focus on the consequences for ${}^{40}\text{Ar}/{}^{39}\text{Ar}$ standards and a few other selected samples shown in Table 5. Although our results in general greatly improve consistency between ages determined by the ${}^{40}\text{Ar}/{}^{39}\text{Ar}$ system and those determined by other techniques, in some cases the new calibrations reveal inconsistencies that warrant further consideration. The recalculated age of FCs – 28.305 \pm 0.036 Ma, including systematic errors – is an iconic manifestation of the results. For the age uncertainty calculation in this case we assumed a relative uncertainty of 0.11% for the R-value, corresponding to the average reproducibility of ${}^{40}\text{Ar}*/{}^{39}\text{Ar}_{K}$ reported by Renne et al. (1998). This age is distinguishable from the astronomically-tuned age of 28.201 ± 0.023 Ma (Kuiper et al., 2008) at 95% confidence, and this discrepancy remains to be further investigated. Recalculated ⁴⁰Ar/³⁹Ar ages of several other standards that have been intercalibrated with FCs are shown in Table 5. Alder Creek sanidine (ACs), Taylor Creek sanidine (TCs), GA-1550 biotite, and Hb3gr hornblende all yield age accuracies better than or equal to 0.16% when systematic errors in λ_{ϵ} , λ_{β} , and κ_{FCs} are included. NL-25 hornblende has a larger uncertainty of 0.38% including systematic errors, largely due to the imprecision (0.72%) of the *R*-value relating it to FCs.

The recalculated ${}^{40}\text{Ar}/{}^{39}\text{Ar}$ age $(1.206 \pm 0.002 \text{ Ma})$ including systematic errors) of the Alder Creek sanidine (ACs) standard may be compared with independent estimates for the age of the Cobb Mountain geomagnetic polarity subchron recorded by its host rhyolite. Unfortunately the various estimates for the age of this event are too dispersed to provide an authoritative independent validation (e.g., see discussion by Nomade et al., 2005 and references therein). However, we note that the most accurate age range (1.215-1.190 Ma) is likely that inferred by Channell et al. (2002) from high sedimentation rate cores in the Iceland Basin, where high resolution magnetostratigraphy and oxygen isotope stratigraphy permitted astronomical tuning with an estimated resolution of ± 5 ka. Thus the recalculated ⁴⁰Ar/³⁹Ar age for ACs remains consistent with this rather broad independently estimated age range.

The recalculated 40 Ar/ 39 Ar age (99.77 ± 0.11 Ma including systematic errors) for GA-1550 biotite is significantly older than the 99.12 ± 0.07 Ma age inferred from 206 Pb/ 238 U data for 10 of 20 zircons analyzed from sample RSES01-98 by Schoene et al. (2006). These ages are clearly distinct at 95% confidence, and an older age for Ar closure in biotite than for Pb retention in zircon is unexpected for this rock unit, which is an intrusive monzonite. One of

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Standard or datum	R	$\pm \sigma$	$\pm \sigma$	Age (Ma)	$\pm \sigma$ (Ma)	±σ (%)
BTs	0.02729	0.00013	0.48	0.7784	0.0037	0.48
ACs	0.04229	0.00006	0.14	1.2061	0.0019	0.16
FCs	1.0000	0.0011	0.11	28.305	0.036	0.13
TCs	1.0112	0.0010	0.10	28.619	0.034	0.12
KTB	2.3650	0.0015	0.06	66.236	0.060	0.09
GA-1550	3.5958	0.0031	0.09	99.769	0.108	0.11
РТВ	9.4918	0.0038	0.04	252.27	0.18	0.07
Hb3gr	51.878	0.0592	0.11	1080.4	1.1	0.10
NL-25	211.42	1.52	0.72	2648.1	10.1	0.38

R-values are calculated using Eq. (2) and the following sources: BTs (Bishop tuff)- Sarna-Wojcicki et al. (2000); ACs – Nomade et al. (2005); TCs – Renne et al. (1998); KTB (Cretaceous/Tertiary boundary) – Swisher et al. (1993); GA –1550-Jourdan and Renne (2007); PTB (Permian/Triassic boundary) – Renne et al. (1995); Hb3gr – Jourdan and Renne (2007); NL-25 – Jourdan and Renne (2007). Age errors were calculated using Monte Carlo simulation (EA-3).

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the twenty zircons analyzed by Schoene et al. (2006) yielded a $^{206}\text{Pb}/^{238}\text{U}$ age of 99.64 \pm 0.03 Ma, indistinguishable from the $^{40}\text{Ar}/^{39}\text{Ar}$ age, raising the possibility that the younger cluster at 99.12 \pm 0.07 Ma is biased by Pb-loss. This is speculative, however, and we cannot confidently explain the discrepancy.

The age of the Cretaceous/Tertiary boundary, as represented by ⁴⁰Ar/³⁹Ar data for sanidine from the IRZ tonstein dated by Swisher et al. (1993), is recalculated to 66.236 ± 0.060 Ma, including systematic errors. This age is slightly older than that of 65.957 ± 0.040 Ma (Kuiper et al., 2008) estimated from astronomical tuning, and the difference (279 \pm 72 ka) is significant at the 95% confidence level. The disagreement in this case is not simply a further manifestation of previously noted inconsistency between our results and those of astronomical tuning in the Miocene Melilla basin, because the latter was derived from completely independent analysis of strata at Zumaia, Spain. While the difference is not distinguishable at 95% confidence from that of a 405 ka eccentricity cycle (the basis of the orbital tuning age), we also note that the original source of the ⁴⁰Ar/³⁹Ar age we recalculated (Swisher et al., 1993) did not document such details as neutron fluence gradients to justify an extremely precise (0.05% relative precision) Jvalue (a measure of the neutron fluence) ascribed to the sample. Hence it is unclear whether the ⁴⁰Ar/³⁹Ar age is as precise as stated. In any case, the discrepancy is worthy of further investigation.

The recalculated ${}^{40}\text{Ar}/{}^{39}\text{Ar}$ age (252.3 \pm 0.2 Ma, including systematic errors) of the Permian–Triassic boundary involves one of the data pairs (Dt3) used as a constraint in this study. The ${}^{40}\text{Ar}/{}^{39}\text{Ar}$ age is indistinguishable from the ${}^{206}\text{Pb}/{}^{238}\text{U}$ zircon age of 252.4 \pm 0.2 Ma reported by Mundil et al. (2004). Making an *ad hoc* residence time correction as discussed previously renders a ${}^{206}\text{Pb}/{}^{238}\text{U}$ age of 252.3 \pm 0.2 Ma, which also coincides with the best-fit age – the age that is most consistent with all other data – given in Table 4c.

We recalculated the ⁴⁰Ar/³⁹Ar age of the Bishop Tuff from data originally reported by Sarna-Wojcicki et al. (2000) in order to compare the result with a ${}^{238}U/{}^{206}Pb$ zircon age reported by Crowley et al. (2007). These results fail our criteria for use in the regression employed here because of pronounced discordance between 238 U/ 206 Pb and 235 U/ 207 Pb ages, attributed by Crowley et al. (2007) to initial Pa/U disequilibrium. Seventeen of nineteen zircons reported by Crowley et al. (2007) yielded a ²³⁸U/²⁰⁶Pb age of 767.1 ± 0.5 ka, indistinguishable (at 95% confidence) from an inferred eruption age (770.4 \pm 1.8 ka, recalculated from the ⁴⁰Ar/³⁹Ar data of Sarna-Wojcicki et al., 2000, using the standard calibration of Renne et al., 1998). Because of this inferred agreement between ${}^{238}U/{}^{206}Pb$ and ⁴⁰Ar/³⁹Ar ages, Crowley et al. (2007) concluded that "high-precision ID-TIMS geochronology can resolve magma chamber dynamics of <1 Ma silicic eruptions at the millennial scale". Moreover, because the majority of zircons were inferred to yield tightly clustered ²³⁸U/²⁰⁶Pb ages (after initial Th/U disequilibrium correction) coincident with eruption, Crowley et al. (2007) questioned the validity of ion probe U-Pb dates from Bishop Tuff zircons which have been inferred to document substantial residence times Table 4c

Optimization results: optimal values of other parameters estimated in the regression analysis. As discussed in the text, the optimization scheme yields refined values for the measurements of *R* (denoted R_i^*); these, with the optimal values of the other parameters κ_{FCs} , λ_e and λ_β , imply refined values for τ (denoted τ_i^*).

Sample	R_i^*	$ au_i^*$
		(Ma)
79CE	6.727×10^{-5}	1.919×10^{-3}
IP	0.6076	17.251
P1T-2	4.5531	125.43
PR94-7	4.8319	132.84
NMB	7.4616	201.21
MSG-09	9.0676	241.73
SH-10	9.4935	252.33
D3T	9.4919	252.29
SH-09	9.5087	252.71
PD97-2	9.5016	252.53
SH-27	9.5289	253.21
SH-16	9.5408	253.50
SH-08	9.6975	257.38
PaV	12.604	327.88
DK-LQ	18.114	454.37
F-239	52.796	1094.8
EGB-032	135.72	2066.9

and protracted crystallization intervals (Simon and Reid, 2005). However, our recalculated 40 Ar/ 39 Ar age of the Bishop Tuff is 778 ± 4 ka (systematic errors included), resolvably older (by 11 ± 4 ka) than the age inferred from 238 U/ 206 Pb data by Crowley et al. (2007), suggesting that either the 40 Ar/ 39 Ar system in sanidine somehow records a pre-eruptive event, or that the 238 U/ 206 Pb age reported by Crowley et al. (2007) is biased for some reason (e.g., an inaccurate correction for initial Th/U disequilibrium).

A further implication of the calibration presented here is that the Brunhes/Matuyama geomagnetic polarity transition, which predates the eruption of the Bishop Tuff (Sarna-Wojcicki et al., 2000 and references therein), is older than 778 \pm 4 ka. This constraint is consistent with the most precise available ⁴⁰Ar/³⁹Ar age for transitional polarity lavas from Maui ascribed to the transition (Coe et al., 2004), which is recalculated with the constants derived herein to 784 ± 2 ka, and is also consistent with some age estimates for the transition derived from orbital tuning, such as that (781 ka) summarized by Lourens et al. (2004). At face value, these results are marginally inconsistent with some other orbitally tuned chronologies for the Brunhes/Matuyama transition, such as those (~773 ka) reported by Channell et al. (2004). These apparent inconsistencies should be the focus of future investigations.

6. CONCLUSIONS

Here we have found values for the ⁴⁰K decay constants (λ_{ε} and λ_{β}), and the ⁴⁰Ar*/⁴⁰K ratio (κ_{FCs}) of the FCs standard that best fit: (i) constraints from activity-based measurements of the ⁴⁰K decay constants, (ii) independent measurements of the age of the FCs standard, and (iii) paired U–Pb and ⁴⁰Ar/³⁹Ar ages on carefully selected samples. First, these esti-

mates of λ_{ε} and λ_{β} are significantly more accurate than obtainable from existing disintegration counting data alone. Second, uncertainties in these three parameter estimates are correlated, and the correlations must be accounted for in computing uncertainties of ⁴⁰Ar/³⁹Ar ages based on these results. Third, the fitting exercise shows that all these diverse constraints used as input data are mutually consistent at their stated uncertainties. Fourth, we have in effect calibrated the ⁴⁰Ar/³⁹Ar system with existing measurements of both the 40 K and 238 U decay constants. Thus, to the extent that uncertainties in these parameters represent the principal sources of systematic error in ⁴⁰Ar/³⁹Ar ages, our results imply approximately an order of magnitude improvement in the accuracy of the ⁴⁰Ar/³⁹Ar method relative to previous calibrations, over most of geologic time. Consequently, the ⁴⁰Ar/³⁹Ar method is now capable of accuracy comparable to that of the U/Pb system.

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APPENDIX A

A.1. Sample notes

Most of the data used herein have previously been published, and the reader is referred to the original reference for information about the samples and data. Samples for which ⁴⁰Ar/³⁹Ar and/or ²⁰⁶Pb/²³⁸U data are being presented for the first time are discussed in further detail. In addition, notes about new data and other features of possible interest are included by sample.

79*CE.* 40 Ar/ 39 Ar (sanidine) data from this sample of pumice from the Plinian eruption of Vesuvius volcano near Pompei, Italy, were originally reported by Renne et al. (1997). New results from an additional sample were reported by Renne and Min (1998), whose combined results were corrected for the year between the two sets of analyses. The calendar year age is that in 1998 when the more recent isotopic data were measured.

IP. Sample IP is a rhyolite ignimbrite from the Gyulakeszi Rhyolite Tuff Formation exposed near Ipolytarnóc, Hungary, where it is associated with well-preserved Miocene ichnofossils. Both U/Pb and 40 Ar/ 39 Ar data were reported by Palfy et al. (2007) and those results are used herein.

PIT-2. (41° 35.6' N, 120° 35.7' E) Sample P1T-2 is a rhyolite tuff intercalated within fossil bearing Early Cretaceous terrestrial sediments of the Yixian Fm. from the Sihetun locality. Swisher et al. (1999) reported a weighted mean $^{40}\mathrm{Ar}/^{39}\mathrm{Ar}$ age of 124.60 \pm 0.25 Ma for 35 single crystals of sanidine, and a plateau age of 124.6 ± 0.1 Ma for a stepheated multigrain aliquot. The sample is equivalent to, though not from the same collection as, sample YX07-4 from Chang et al. (2009), who reported a weighted mean age of 123.8 ± 0.2 Ma for 17 single crystals of sanidine, and plateau ages of 124.1 ± 0.2 and 124.1 ± 0.2 Ma for step-heated multigrain aliquots. While the data of Chang et al. (2009) are more consistent with the other ⁴⁰Ar/³⁹Ar-²⁰⁶Pb/²³⁸U data pairs than those of Swisher et al. (1999), there is no objective basis for excluding either data set a priori, except that Swisher et al. (1999) did not report relative isotope abundances corresponding to their plateau age, nor any information about the basis for the J-value or its contribution to the uncertainty budget. Accordingly, the data corresponding to the mean plateau age $(124.1 \pm 0.16 \text{ Ma})$ of Chang et al. (2009) are used for the present study. New U/Pb analyses of 6 annealed and chemically abraded zircon crystals from P1T-2 yielded a weighted mean ${}^{206}\text{Pb}/{}^{238}\text{U}$ age of 125.65 \pm 0.17 Ma (MSWD 0.63).

PR94-7. This rock is a quartz-plagioclase-sanidine phyric rhyolite ash-flow tuff from the Arequita Fm., 4 km west of Lascano in southeastern Uruguay, where it appears to represent a late stage eruption of the voluminous Paraná-Etendeka flood volcanic province. 206Pb/238U (zircon) and ⁴⁰Ar/³⁹Ar data (sanidine) were obtained for this study. The ignimbrite is approximately 50 m thick, and is moderately to densely welded near its base. The sample was collected from ~ 20 m above the base. 16 individual zircons crystals from sample PR94-7 were subjected to either air abrasion or pretreated in ultrasonically agitated, hot HF aimed at minimizing the effects of Pb loss. 13 analyses yield a ${}^{206}\text{Pb}/{}^{238}\text{U}$ age of 133.06 ± 0.26 Ma (MSWD 1.0). An additional 6 zircons show distinctly older and/or discordant ages suggesting inheritance. Optically clear sanidine from the 20 to 40 mesh fraction (417-833 µm) of PR94-7 was separated by conventional methods and irradiated for 25 h in the OSTR reactor along with Fish Canyon sanidine (FCs) in three bracketing sample positions in an irradiation disk as described previously. Five single-crystal and one tencrystal sample were analyzed by incremental heating to fusion using a CO₂ laser system with integrator lens. The $^{40}\mathrm{Ar}^{*/^{39}\!Ar_{K}}$ value attributed to the standard (Table 3) is the arithmetic mean of three closely bracketing FCs aliquots. All of the experiments yielded well-defined age plateaux in mutual agreement within analytical errors, yielding MSWD = 0.5. Lower precision of the single-crystal experiments is due to the larger influence of background corrections on smaller argon yields.

NMB. The North Mountain basalt is part of the Central Atlantic Magmatic Province (CAMP). ²⁰⁶Pb/²³⁸U zircon data were reported by Schoene et al. (2006) and ⁴⁰Ar/³⁹Ar on plagioclase were reported by Jourdan et al. (2009). We use these data as reported in the original sources.

MSG-09. (45° 54.6′ N, 8° 54.6′ E) This bentonite occurs at the Anisian/Ladinian Grenzbitumenzone near Monte San Giorgio, Switzerland. 206 Pb/ 238 U (zircon) data and 40 Ar/ 39 Ar (sanidine) data from this bentonite were reported by Mundil et al. (2010).

SH-10. The sample is a bentonite equivalent to GS-1, for which 40 Ar/ 39 Ar on plagioclase were reported by Renne et al. (1995). New 40 Ar/ 39 Ar (plagioclase) data were obtained for this study by step-heating 3 multigrained aliquots that yielded indistinguishable plateau ages. Individual analysis of 13 crystals, combined with 38 previously reported by Renne et al. (1995), fail to reveal any xenocrystic contamination. The new results are combined with previous data as shown in Table 3. The FCs grains were loaded in the same well as the sample, in an irradiation disk similar to those depicted by Renne et al. (1998).

D3T. ²⁰⁶Pb/²³⁸U zircon data from this bentonite were reported by Mundil et al. (2004). This sample is equivalent to sample C-2, for which Renne et al. (1995) reported ⁴⁰Ar/³⁹Ar (sanidine) data. The data used herein are derived from these two sources.

SH-09. ²⁰⁶Pb/²³⁸U zircon data from this bentonite were reported by Mundil et al. (2004). New ⁴⁰Ar/³⁹Ar data (sanidine) were obtained for this study. Two multigrained aliquots were heated incrementally and yielded indistinguishable plateau ages, producing the data used herein (Table 3). Ten single crystal analyses yielded no evidence of xenocrysts. The FCs grains were loaded in the same well as the sample, in an irradiation disk similar to those depicted by Renne et al. (1998).

PD97-2. ⁴⁰Ar/³⁹Ar data (biotite) from this bentonite in the lower Quartermaster Fm. in Palo Duro State Park, Texas, were reported by Renne et al. (2001). New U–Pb data were obtained for this study. 7 out of 7 HF-leached zircons yielded a ²⁰⁶Pb/²³⁸U age of 253.02 \pm 0.50 Ma (MSWD 1.2).

SH-27. 206 Pb/ 238 U zircon data from this bentonite were reported by Mundil et al. (2004). New 40 Ar/ 39 Ar data (sanidine) were obtained for this study. The sanidine crystals were sufficiently large that incremental heating of individual grains was possible, and nine such grains yielded indistinguishable plateau ages with MSWD of 0.44. The FCs grains were loaded in the same well as the sample, in an irradiation disk similar to those depicted by Renne et al. (1998).

SH-16. ²⁰⁶Pb/²³⁸U zircon data from this bentonite were reported by Mundil et al. (2004). New ⁴⁰Ar/³⁹Ar data (sanidine) were obtained for this study. A single multigrained aliquot was analyzed and yielded an age plateau comprising 100% of the ³⁹Ar released. Individual total fusion of 28 grains revealed no evidence of xenocrysts. The FCs grains were loaded in the same well as the sample, in an irradiation disk similar to those depicted by Renne et al. (1998).

SH-08. ²⁰⁶Pb/²³⁸U zircon data from this bentonite were reported by Mundil et al. (2004). New ⁴⁰Ar/³⁹Ar data (sanidine) were obtained for this study. Four multigrained aliquots were incrementally heated and yielded indistinguishable plateau ages with MSWD = 0.76. Individual total fusion of 23 grains revealed no evidence of xenocrysts. The FCs grains were loaded in the same well as the sample, in an irradiation disk similar to those depicted by Renne et al. (1998).

PaV. (36° 36.1′ S, 151° 36.9′ E) Sample PaV is from a ca. 100 m thick volcanic unit with quartz latite (toscanite) composition of Lower Carboniferous age exposed in a railway cut near Paterson (NSW, Australia) which is bracketed by clastic terrestrial sediments and used as a marker bed for the onset of the Kiaman magnetic superchron in Australia as well as the late Paleozoic Gondwana glaciation (Irving, 1966; Claoué-Long et al., 1995). U/Pb SHRIMP and ID-TIMS analyses by Claoué-Long et al. (1995) yielded ages of 330.1 ± 3.0 and 328.9 ± 4.2 Ma from the former and 327.7 ± 2.0 Ma from the latter technique. New 206 Pb/ 238 U (zircon) and ⁴⁰Ar/³⁹Ar (hornblende) data were obtained for this study. U/Pb analyses of annealed and chemically abraded zircons yield a 206 Pb/ 238 U age of 328.0 ± 0.5 Ma (MSWD 0.6) from six out of six crystals, in agreement with but more precise than the U/Pb ages from Claoué-Long et al. (1995). Five multigrain hornblende aliquots yielded 40 Ar/ 39 Ar plateaux with >90% of the 39 Ar released, yielding a weighted mean ${}^{40}\text{Ar}*/{}^{39}\text{Ar}_{K}$ of 37.303 ± 0.032 , MSWD = 2.0 and probability of fit = 0.094. To account for the excess scatter, the uncertainty was expanded by Students $t \times \text{SORT}(\text{MSWD})$ to give an uncertainty of $\sigma_{\rm R} = 0.065$. The FCs grains were loaded in the same well as the sample, in an irradiation disk similar to those depicted by Renne et al. (1998).

DK-LQ. (37° 55.6′ N, 84° 55.9′ W). This sample is the Deicke bentonite collected at Lexington quarry, Kentucky. ⁴⁰Ar/³⁹Ar data for biotite single crystals were reported by Min et al. (2001), who employed a strict concordance criterion (plateaux comprising 100% of ³⁹Ar released; met by 9 of 11 crystals) to reject samples suspected of showing ³⁹Ar recoil redistribution effects from interlayered chlorite alteration present in some grains. New U/Pb zircon data were obtained for this study. Sixteen individual zircons were annealed and chemically abraded. Ten yielded a concordant cluster with a ²⁰⁶Pb/²³⁸U age of 454.59 ± 0.56 Ma (MSWD 1.8); five are distinctly older and clearly represent inherited grains; and 1 is resolvably younger and is inferred to reflect Pb-loss.

F-239. ²⁰⁷Pb/²⁰⁶Pb (zircon) and ⁴⁰Ar/³⁹Ar (anorthoclase) ages were reported by Min et al. (2000) for the Palisade Rhyolite, a densely welded ignimbrite of the North Shore volcanic group in Minnesota. Schoene et al. (2006) reported U/Pb (zircon) data for the same unit (sample #MS99-30), and their result is used herein. The ⁴⁰Ar/³⁹Ar results of Min et al. (2000) are adopted.

EGB-032. This sample is an unmetamorphosed hornblende–biotite dacite porphyry from the Eglab region of the Requibath massif, west Africa. U/Pb zircon data from this sample were reported by Schoene et al. (2006). New 40 Ar/ 39 Ar data (hornblende) were obtained from eleven single crystals, each yielding a plateau comprising >90% of the 39 Ar released, for this study. The individual plateau values of 40 Ar*/ 39 Ar_K (*F*-values of Table 3) yield MSWD = 0.75. The FCs grains were loaded in the same well as the sample, in an irradiation disk similar to those depicted by Renne et al. (1998).

A.2. Calculation of ages and uncertainties

The ⁴⁰Ar/³⁹Ar age of an unknown based on the results presented herein can be calculated by solving Eq. (3) for τ_I , changing subscripts as appropriate, and substituting t_{unk} for τ_I .

$$t_{\text{unk}} = \frac{1}{(\lambda_{\text{tot}})} \ln \left[1 + \left(\frac{\lambda_{\text{tot}}}{\lambda_{\varepsilon}} \right) R_{\text{FCs}}^{\text{unk}} \cdot \kappa_{\text{FCs}} \right]$$
(A.2.1)

If the unknown is run with the FCs standard, R is as defined by Eq. (2). If another standard is used, the ${}^{40}\text{Ar}*/{}^{39}\text{Ar}_{\text{K}}$ of the unknown must be referenced to FCs using the relationship:

$$R_{\text{FCs}}^{\text{unk}} \equiv R_{\text{std1}}^{\text{unk}} \cdot R_{\text{std2}}^{\text{std1}} \cdot \dots \cdot R_{\text{FCs}}^{\text{stdn}}$$
(A.2.2)

and appropriate intercalibration factors as given by, e.g., Renne et al. (1998) for several widely used standards. Alternatively, *R* values relating the unknown to another standard may be used provided that a value of κ appropriate to that standard is also used. K–Ar ages may be simply computed by substituting (i) the measured ⁴⁰Ar*/³⁹Ar_K for κ_{FCs} and (ii) a value of 1 for R_{FCs}^{unk} in Eq. (A.2.1).

The linear error propagation formula given by Eq. (7) makes use of the following partial derivatives:

$$\frac{\partial t}{\partial \lambda_{\varepsilon}} = \frac{-1}{(\lambda_{\varepsilon} + \lambda_{\beta})} \left[t + \frac{\lambda_{\beta} \kappa_{\text{FCs}} R}{(\lambda_{\varepsilon}^2 \kappa_{\text{FCs}} R + \lambda_{\varepsilon} \lambda_{\beta} \kappa_{\text{FCs}} R + \lambda_{\varepsilon}^2)} \right] \quad (A.2.3)$$

$$\frac{\partial t}{\partial \lambda_{\beta}} = \frac{1}{(\lambda_{\varepsilon} + \lambda_{\beta})} \left[\frac{\kappa_{\text{FCs}} R}{(\lambda_{\varepsilon} \kappa_{\text{FCs}} R + \lambda_{\beta} \kappa_{\text{FCs}} R + \lambda_{\varepsilon})} - t \right]$$
(A.2.4)

$$\frac{\partial t}{\partial \kappa_{\rm FCs}} = \frac{R}{(\lambda_{\varepsilon} \kappa_{\rm FCs} R + \lambda_{\beta} \kappa_{\rm FCs} R + \lambda_{\varepsilon})} \tag{A.2.5}$$

$$\frac{\partial t}{\partial R} = \frac{\kappa_{\rm FCs}}{(\lambda_{\varepsilon}\kappa_{\rm FCs}R + \lambda_{\beta}\kappa_{\rm FCs}R + \lambda_{\varepsilon})} \tag{A.2.6}$$

and the covariances given in Table 4b.

To illustrate the application of these equations, we detail the recalculation of the age of the Cretaceous/Tertiary boundary as given in Section 5.3. This is based on an age of 65.16 \pm 0.04 Ma for the IrZ coal, as originally reported by Swisher et al. (1993) relative to an age of 27.84 Ma for FCs. We first calculate $R_{\rm FCs}^{\rm IrZ}$ using Eq. (2), using the value of $\lambda_{\rm tot}$ (5.543 \times 10⁻¹⁰ a⁻¹) used in the original reference:

$$R_{\text{FCs}}^{\text{IrZ}} \equiv \frac{\left({}^{40}Ar^{*}\right)^{39}Ar_{K}\right)_{\text{IrZ}}}{\left({}^{40}Ar^{*}\right)^{39}Ar_{K}\right)_{\text{FCs}}} \equiv \frac{\left(e^{\lambda_{\text{tot}}I_{\text{IrZ}}} - 1\right)}{\left(e^{\lambda_{\text{tot}}I_{\text{FCs}}} - 1\right)}$$
$$= \frac{\left(e^{5.543 \times 10^{-10} \cdot 6.516 \times 10^{7}} - 1\right)}{\left(e^{5.543 \times 10^{-10} \cdot 2.784 \times 10^{7}} - 1\right)} = 2.36496$$
(A.2.7)

Application of Eq. (A.2.1) and the data in Table 4a yields an age of 66.236 Ma.

Uncertainty arising from the decay constants was neglected in the reported age, therefore by linear error propagation, the uncertainty in the value of R_{FCs}^{IrZ} is given by:

$$\sigma_{R_{\rm IrZ}} \equiv \frac{\lambda_{\rm tot} e^{\lambda_{\rm tot} \cdot t_{\rm IrZ}}}{(e^{\lambda_{\rm tot} \cdot t_{\rm FCs}} - 1)} \sigma_{t_{\rm IrZ}} = 0.00148 \text{ Ma}$$
(A.2.8)

Now the values of the partial derivatives can be computed from Eqs. A.2.3, A.2.4, A.2.5, A.2.6:

$$\frac{\partial t}{\partial \lambda_{\varepsilon}} = -1.1322 \times 10^{18} a^2 \tag{A.2.3}$$

$$\frac{\partial t}{\partial \lambda_{\beta}} = -2.1670 \times 10^{15} a^2 \tag{A.2.4}$$

$$\frac{\partial t}{\partial \kappa_{\rm FCs}} = 3.9612 \times 10^{10} \ a \tag{A.2.5}$$

$$\frac{\partial t}{\partial R} = 2.7499 \times 10^7 \ a \tag{A.2.6}$$

Application of these values, and the covariances given in Table 4b, to Eq. (7) yields an age uncertainty of $\sigma_{r_{\rm IrZ}} = 0.041$ Ma. As noted in the text however, linear error propagation is somewhat inaccurate for very precise R-values such as in this case, where $\sigma_{R_{\rm IrZ}}/R_{\rm IrZ} = 0.06\%$. Instead, Monte Carlo simulation is preferred and in this case yields an uncertainty of ± 0.060 Ma as shown in Table 5, based on the distributions given in EA-3.

An Excel spreadsheet which performs age calculations and both linear and Monte Carlo error estimation, based on the results presented herein, is available on request from the corresponding author (P.R.R.).

APPENDIX B. SUPPLEMENTARY DATA

Supplementary data associated with this article can be found, in the online version, at doi:10.1016/j.gca.2010.06.017.

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